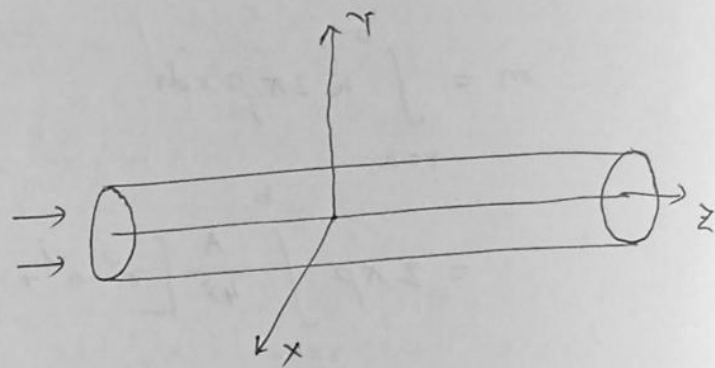


Case (III):

Steady flow in pipes of elliptic cross-section:

Let us consider the incompressible viscous fluid through a pipe having elliptic cross-section.



Also let z-axis be taken in the direction of flow along the axis of the pipe.

In steady case, the equation of continuity and motion are.

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \nu \nabla^2 u$$

$$u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial y} + \nu \nabla^2 v$$

$$\& u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial z} + \nu \nabla^2 w$$

$$\text{where } \nabla^2 \equiv \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

Here, $u = v = 0, w \neq 0$.

therefore the above equation reduce to

$$\frac{\partial w}{\partial z} = 0 \rightarrow \textcircled{I}$$

$$-\frac{1}{\rho} \frac{\partial p}{\partial x} = 0 \rightarrow \textcircled{II}$$

$$-\frac{1}{\rho} \frac{\partial p}{\partial y} = 0 \rightarrow \textcircled{III}$$

$$-\frac{1}{\rho} \frac{\partial p}{\partial z} + \nu \left[\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} \right] = 0 \rightarrow \textcircled{IV}$$

from \textcircled{II} & \textcircled{III} , p is independent of x and y only,

therefore $p = p(z)$.

Now, the equation (iv) is solvable if each term is constant

$$\text{i.e. } -\frac{1}{\rho} \frac{\partial p}{\partial z} = A \text{ (constant)}$$

$$\therefore \gamma \left[\frac{\partial \tilde{w}}{\partial x^2} + \frac{\partial \tilde{w}}{\partial y^2} \right] + A = 0$$

$$\Rightarrow \frac{\partial \tilde{w}}{\partial x^2} + \frac{\partial \tilde{w}}{\partial y^2} = -\frac{A}{\gamma} \rightarrow \text{(vi)}$$

This equation of w is to be solved under the boundary conditions

$$w = 0 \text{ at } \frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

$$\text{and } w \text{ is continuous in } \frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

We take the trial solution of w as

$$w = B \left(1 - \frac{x^2}{a^2} - \frac{y^2}{b^2} \right) \rightarrow \text{(vii)}$$

This satisfies the boundary conditions and we fit it in the differential equation to find the constant B .

Putting this value of w in (vi), we get,

$$B \left(-\frac{2}{a^2} \right) + B \left(-\frac{2}{b^2} \right) = -\frac{A}{\gamma}$$

$$\Rightarrow 2B \cdot \frac{a^2 + b^2}{a^2 \times b^2} = \frac{A}{\gamma}$$

$$\Rightarrow B = \frac{A}{2\gamma} \left(\frac{a^2 b^2}{a^2 + b^2} \right)$$

Substituting in (vii),

$$w = \frac{A}{2\gamma} \left(\frac{a^2 b^2}{a^2 + b^2} \right) \left(1 - \frac{x^2}{a^2} - \frac{y^2}{b^2} \right)$$

Now, the flux of the fluid over the area of the

ellipse is given by

$$\iint W \, dxdy$$

$$= \iint \frac{A}{2\gamma} \frac{a^{\sqrt{x}} b^{\sqrt{y}}}{a^{\sqrt{x}+b^{\sqrt{y}}} \left(1 - \frac{\sqrt{x}}{a} - \frac{\sqrt{y}}{b}\right)} \, dxdy,$$

where y varies from $-\frac{b}{a} \sqrt{a-x}$ to $\frac{b}{a} \sqrt{a-x}$ and
 x varies from $-a$ to a .

Now,

$$\iint W \, dxdy$$

$$= \frac{A}{2\gamma} \frac{a^{\sqrt{x}} b^{\sqrt{y}}}{a^{\sqrt{x}+b^{\sqrt{y}}}} \left[\iint dxdy - \frac{1}{a} \iint \sqrt{x} \, dxdy - \frac{1}{b} \iint \sqrt{y} \, dxdy \right]$$

Now, $a \frac{b}{a} \sqrt{a-x}$

$$\int_{x=-a}^a \int_{y=-\frac{b}{a} \sqrt{a-x}}^{\frac{b}{a} \sqrt{a-x}} dxdy = \int_{-a}^a \left[y \right]_{-\frac{b}{a} \sqrt{a-x}}^{\frac{b}{a} \sqrt{a-x}} dx$$

$$= \int_{-a}^a \frac{2b}{a} \sqrt{a-x} \, dx$$

$$= \frac{2b}{a} \int_{\pi}^0 \sqrt{a-a \cos \theta} (-a \sin \theta) d\theta \quad \left| \begin{array}{l} \text{Putting} \\ x = a \cos \theta \\ dx = -a \sin \theta d\theta \end{array} \right.$$

$$= 2ab \frac{1}{2} \int_0^{\pi} (1 - \cos 2\theta) d\theta$$

$$= ab \left[\theta - \frac{\sin 2\theta}{2} \right]_0^{\pi}$$

$$= ab [\pi - 0]$$

$$= \pi ab.$$

$$\iint x^{\nu} dx dy = \int_{-a}^a \int_{-\frac{b}{a}\sqrt{a^2-x^2}}^{\frac{b}{a}\sqrt{a^2-x^2}} x^{\nu} dx dy$$

$$= \int_{-a}^a x^{\nu} \left[y \right]_{-\frac{b}{a}\sqrt{a^2-x^2}}^{\frac{b}{a}\sqrt{a^2-x^2}} dx$$

$$= \int_{-a}^a \frac{2b}{a} x^{\nu} \sqrt{a^2-x^2} dx$$

$$= \int_{\pi}^0 \frac{2b}{a} \cdot a^{\nu} \cos^{\nu} \theta \cdot a \sin \theta \cdot (-a \sin \theta) d\theta$$

$$= 2a^3 b \int_0^{\frac{\pi}{2}} \cos^{\nu} \theta \sin^2 \theta d\theta$$

$$= 2a^3 b \cdot 2 \int_0^{\frac{\pi}{2}} \cos^{\nu} \theta \sin \theta d\theta$$

$$= 4a^3 b \cdot \frac{\pi}{16}$$

$$= \frac{\pi a^3 b}{4}$$

and $\iint y^{\nu} dx dy = \int_{-a}^a \left[\frac{y^3}{3} \right]_{-\frac{b}{a}\sqrt{a^2-x^2}}^{\frac{b}{a}\sqrt{a^2-x^2}} dx$

$$= \frac{1}{3} \int_{-a}^a \left[\frac{b^3}{a^3} (\sqrt{a^2-x^2})^3 + \frac{b^3}{a^3} (\sqrt{a^2-x^2})^3 \right] dx$$

$$= \frac{2b^3}{3a^3} \int_{-a}^a (\sqrt{a^2-x^2})^3 dx$$

$$= \frac{2b^3}{3a^3} \int_{\pi}^0 a^3 \sin^3 \theta \cdot (-a \sin \theta) d\theta$$

$$= \frac{2b^3 a}{3} \int_0^{\frac{\pi}{2}} \sin^4 \theta d\theta = \frac{2ab^3}{3} \cdot 2 \int_0^{\frac{\pi}{2}} \sin^4 \theta d\theta = \frac{2ab^3}{3} \cdot 2 \cdot \frac{3}{4} \cdot \frac{1}{2} \cdot \frac{\pi}{2} = \frac{\pi ab^3}{4}$$

Hence, $a \frac{b}{a} \sqrt{a^2 - r^2}$

$$\begin{aligned} \iint_{-a - \frac{b}{a} \sqrt{a^2 - r^2}}^{a \frac{b}{a} \sqrt{a^2 - r^2}} W \, dx \, dy &= \frac{A}{2\gamma} \frac{a^2 b^2}{a^2 + b^2} \left[\pi ab - \frac{1}{a^2} \frac{\pi a^3 b}{4} - \frac{1}{b^2} \frac{\pi ab^3}{4} \right] \\ &= \frac{A}{2\gamma} \frac{a^2 b^2}{a^2 + b^2} \left[\pi ab - \frac{\pi ab}{4} - \frac{\pi ab}{4} \right] \\ &= \frac{A\pi}{4\gamma} \frac{a^3 b^3}{a^2 + b^2} \end{aligned}$$

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Steady rotation about the axis:

Let us consider the steady rotation of the incompressible viscous fluid about the axis of z . Suppose

u, v, w be the velocity components along r, θ, z increasing direction. we have considered rotation

only therefore

$$u = w = 0, \quad v \neq 0.$$

Now, the equation of continuity and motion in cylindrical polar coordinates (r, θ, z) are

$$\frac{\partial u}{\partial r} + \frac{u}{r} + \frac{\partial v}{r \partial \theta} + \frac{\partial w}{\partial z} = 0$$

$$u \frac{\partial u}{\partial r} + v \frac{\partial u}{r \partial \theta} + w \frac{\partial u}{\partial z} - \frac{v^2}{r} = -\frac{1}{\rho} \frac{\partial p}{\partial r} + \gamma \left[\nabla^2 u - \frac{2}{r^2} \frac{\partial v}{\partial \theta} - \frac{u}{r^2} \right]$$

$$u \frac{\partial v}{\partial r} + v \frac{\partial v}{r \partial \theta} + w \frac{\partial v}{\partial z} + \frac{uv}{r} = -\frac{1}{\rho} \frac{\partial p}{r \partial \theta} + \gamma \left[\nabla^2 v - \frac{v}{r} + \frac{2}{r^2} \frac{\partial u}{\partial \theta} \right]$$

$$u \frac{\partial w}{\partial r} + v \frac{\partial w}{r \partial \theta} + w \frac{\partial w}{\partial z} = -\frac{1}{\rho} \frac{\partial p}{\partial z} + \gamma \left[\nabla^2 w \right], \text{ where } \nabla^2 \equiv \frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\partial^2}{\partial z^2}$$

Since motion is same for all θ , therefore

$$\frac{\partial p}{\partial \theta} = 0.$$

Now, in this case, the above equation reduce to

$$\frac{\partial U}{\partial \theta} = 0 \rightarrow (i)$$

$$-\frac{v^2}{r} = -\frac{1}{\rho} \frac{\partial p}{\partial r} \rightarrow (ii)$$

$$2) \left[\frac{\partial^2 U}{\partial r^2} + \frac{1}{r} \frac{\partial U}{\partial r} + \frac{\partial^2 U}{\partial z^2} - \frac{v}{rv} \right] = 0 \rightarrow (iii)$$

$$-\frac{1}{\rho} \frac{\partial p}{\partial z} = 0 \rightarrow (iv)$$

Now, $p = p(r)$ only

from (iii) $\frac{d^2 v}{dr^2} + \frac{1}{r} \frac{dv}{dr} - \frac{v}{r^2} = 0 \rightarrow (v)$

[Since motion is in the circle, therefore v is not a function of z .]

If we put, $v = \Omega r$, where Ω is a variable.

$$\frac{dv}{dr} = \Omega + r \frac{d\Omega}{dr}$$

$$\begin{aligned} \frac{d^2 v}{dr^2} &= \frac{d\Omega}{dr} + \frac{d\Omega}{dr} + r \frac{d^2 \Omega}{dr^2} \\ &= 2 \frac{d\Omega}{dr} + r \frac{d^2 \Omega}{dr^2} \end{aligned}$$

Substituting in (v)

$$2 \frac{d\Omega}{dr} + r \frac{d^2 \Omega}{dr^2} + \frac{\Omega}{r} + \frac{d\Omega}{dr} - \frac{\Omega r}{r^2} = 0$$

$$\Rightarrow r \frac{d^2 \Omega}{dr^2} + 3 \frac{d\Omega}{dr} = 0$$

Multiplying both sides by r^2 , we get,

$$r^3 \frac{d^2 \Omega}{dr^2} + 3r^2 \frac{d\Omega}{dr} = 0$$

$$\Rightarrow \frac{d}{dr} \left[r^3 \frac{d\Omega}{dr} \right] = 0$$

$$\Rightarrow r^3 \frac{d\Omega}{dr} = \text{constant} = A \text{ (say)}$$

$$\Rightarrow \frac{d\Omega}{dr} = \frac{A}{r^3}$$

$$\Rightarrow \Omega = -\frac{A}{2r^2} + B$$

$$v = \Omega r = -\frac{A}{2r} + Br. \rightarrow (vi) \quad \#$$