

Drag Formula:

If we impose equal and opposite velocity to the sphere and the liquid together. We find the motion of the sphere of radius 'a' with velocity U through the liquid rest at infinity. Mathematically, the solution to the motion of sphere with velocity U through fluid at rest is obtained. If we remove the 1st term in the expression of ψ

$$\text{Hence } \psi = \frac{3}{4} U a R \left[1 - \frac{1}{3} \frac{a^3}{R^3} \right] \sin^2 \phi \quad \leftarrow \boxed{\text{Drop the 1st term}}$$

$$\therefore v_R = -\frac{3}{2} U \cos \phi \left[\frac{a}{R} - \frac{1}{3} \frac{a^3}{R^3} \right]$$

$$v_\phi = \frac{3}{4} U \sin \phi \left[\frac{a}{R} + \frac{1}{3} \frac{a^3}{R^3} \right]$$

If D is the Drag on the sphere, then work done by the Drag force

= Dissipation of energy

Now, the sphere is moving with velocity U

\therefore Rate of work done by the drag = Dissipation of energy (= DU)

$$= \mu \iiint_V \vec{\tau} \cdot \vec{v} \, dV \quad \left[\begin{array}{l} \text{The fluid is} \\ \text{at rest at} \\ \text{Infinity} \end{array} \right]$$

But $\vec{\tau}$ has only one component

$$\tau_{\phi\theta} = \frac{1}{R} \left[\frac{\partial}{\partial R} (R v_\phi) - \frac{\partial}{\partial \phi} (v_R) \right]$$

$$= \frac{1}{R} \left[\frac{\partial}{\partial R} \left\{ \frac{3}{4} U \sin \phi \left(a + \frac{1}{3} \frac{a^3}{R^3} \right) \right\} - \frac{\partial}{\partial \phi} \left\{ -\frac{3}{2} U \cos \phi \left(\frac{a}{R} - \frac{1}{3} \frac{a^3}{R^3} \right) \right\} \right]$$

$$= \frac{1}{R} \left[\frac{3}{4} U \sin \phi \left(-\frac{1}{3} \frac{2a^3}{R^3} \right) - \left\{ \frac{3}{2} U \sin \phi \left(\frac{a}{R} - \frac{1}{3} \frac{a^3}{R^3} \right) \right\} \right]$$

$$= \frac{1}{R} \left[-\frac{1}{2} \frac{a^3}{R^3} - \frac{3a}{2R} + \frac{1}{2} \frac{a^3}{R^3} \right] U \sin \phi$$

$$= -\frac{3aU \sin \phi}{2R^2}$$

$$DU = \mu \int_{R=a}^{\infty} \int_{\phi=0}^{\pi} \int_{\lambda=0}^{2\pi} \left[\frac{9a^3 U \sin^3 \phi}{4R^4} \cdot R^2 \sin \phi \right] dR d\phi d\lambda$$

$$= 2\pi \mu \int_{R=a}^{\infty} \int_{\phi=0}^{\pi} \frac{9a^3 U \sin^3 \phi}{4R^4} \cdot R^2 \sin \phi d\phi dR$$

$$= 2\pi \mu \frac{9a^3 U}{4} \int_{R=a}^{\infty} \frac{dR}{R^2} \int_{\phi=0}^{\pi} \sin^3 \phi d\phi$$

$$= \frac{9a^3 \pi \mu U}{2} \left[-\frac{1}{R} \right]_a^{\infty} \int_{\phi=0}^{\pi} \sin^3 \phi d\phi$$

$$= \frac{9a^3 U}{2} \mu \pi \frac{4}{3} = 6\pi \mu a U$$

$$\Rightarrow D = 6\pi \mu a U$$

This is known as Drag formula of Stokes.

This is the well known stoke's formula for the drag of a sphere.

We may use this form of drag to find the terminal velocity of a sphere when it moves through the fluid vertically down ward under gravity,

If U be the velocity, ρ' be the density of the sphere, ρ be the density of the fluid outside, then

total force on the sphere = its weight - bouyancy

$$= \frac{4\pi}{3} (\rho' - \rho) a^3 g$$

$$\text{and } \frac{4\pi}{3} (\rho' - \rho) a^3 g = 6\pi\mu a U$$

$$\Rightarrow U = \frac{2}{9\mu} (\rho' - \rho) a^2 g$$

